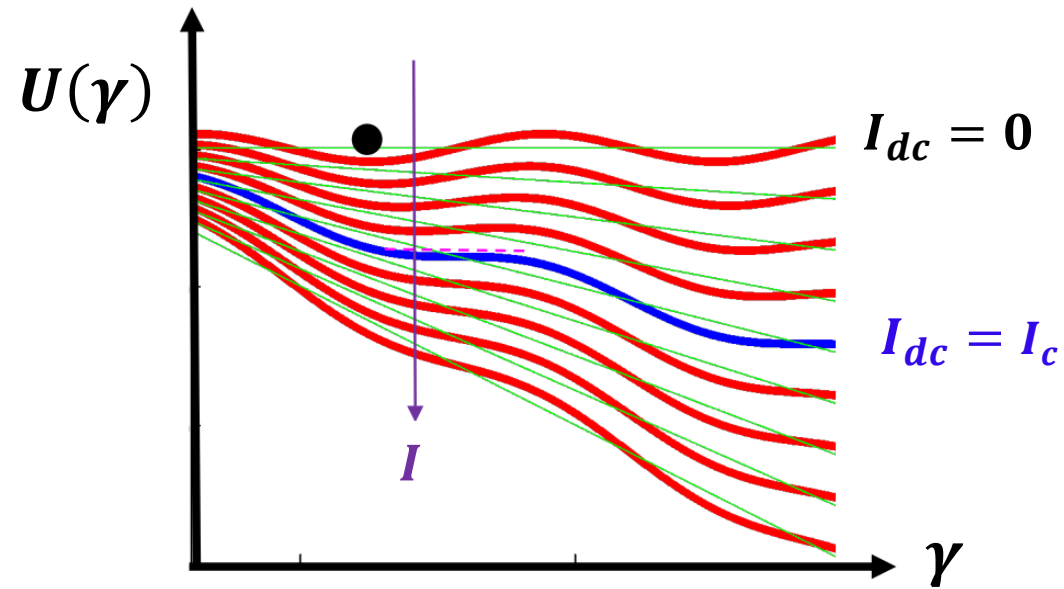
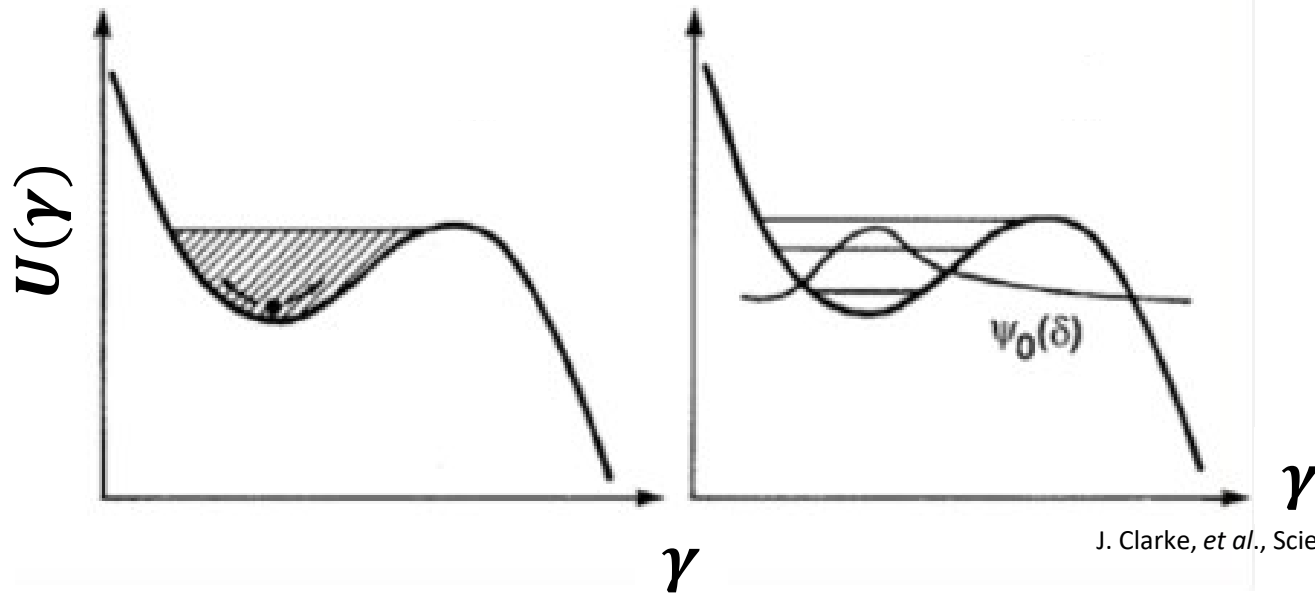


When is a Current-Biased Josephson Junction in the Quantum Limit?

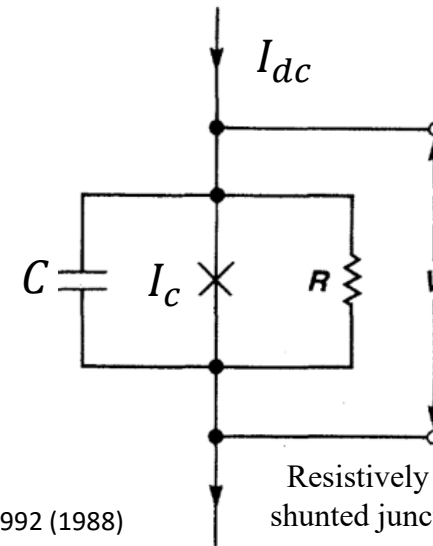
A current-biased Josephson junction has a potential given by $U(\gamma) = -\frac{\Phi_0}{2\pi} (I_{dc} \gamma + I_c \cos \gamma)$, where γ is the gauge-invariant phase difference and Φ_0 is the flux quantum. This creates the tilted washboard potential as a function of gauge-invariant phase.



When is the Problem Quantum as opposed to Classical?



J. Clarke, et al., Science 239, 992 (1988)



DC Josephson equation
 $I = I_c \sin \gamma$

Resistively and capacitively-shunted junction (RCSJ) model

The gauge-invariant phase point oscillates in an (approximately) harmonic oscillator potential characterized by a curvature given by the dc-current-dependent Josephson plasma frequency $\omega_p(I_{dc})$:

$$E_n \approx \hbar \omega_p \left(n + \frac{1}{2} \right) - \text{small anharmonic corrections.} \quad \text{with} \quad \omega_p = \sqrt{\frac{2eI_c}{\hbar C}} \left(1 - \left(\frac{I_b}{I_c} \right)^2 \right)^{1/4}$$

$n = 0, 1, 2, 3, \dots$

To see quantum transitions between the states the thermal energy must be small compared to the energy level spacing

$$k_B T \ll \hbar \omega_p$$

To minimize the effects of loss we also require: $Q = \omega_p RC \gg 1$, in the language of the RCSJ model, shown above

This approach leads to a phase qubit